

1 **Multi-annual monitoring of the water vapor vertical distribution on Mars by** 2 **SPICAM on Mars Express**

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8

9 **Abstract**

10 The distribution of water vapor with altitude has long remained a missing piece of the
11 observational dataset of water vapor on Mars. In this work, we present the first multi-annual survey
12 of water vapor profile covering the altitude range from 0 to 100 km based on the SPICAM/Mars
13 Express occultation measurements. During the aphelion season, water remains confined below 40-
14 60 km for all Martian years observed. The highest altitude where water vapor can be spotted is
15 between 70 and 90 km during the southern summer ($L_s=240-300^\circ$; perihelion season), approaching
16 the transition between the middle and upper atmosphere. In this season, years without a global dust
17 storm (GDS) show a significant moistening of the upper atmosphere (~ 100 ppmv) in the southern
18 hemisphere, confirming a seasonal impact on the hydrogen escape rate. The two observed GDS,
19 in MY28 and MY34, show a substantial disparity in water vapor response. The storm in MY28,
20 which coincides with the southern summer solstice, creates the largest excess of water in both
21 hemispheres at >80 km. This climatology of water vapor will supply a robust statistical basis to
22 address the long-term escape processes of water from Mars.

23

24 **Plain Language Summary**

25 The vertical distribution of water vapor in the Martian atmosphere is key to understanding
26 water transport and its escape from the planet, which in turn helps to explain the fate of water
27 through Mars' history. Recent studies suggest that the transport of water to 80 km can increase the

28 hydrogen escape rate by an order of magnitude and provide evidence for the role of global dust
29 storms (GDS) in the regulation of this process. We monitored the vertical water distribution during
30 multiple Martian years for the first time, including two years with global dust storms. We confirm
31 the efficiency of GDS in delivering water to 80 km and emphasize the role of a regular perihelion
32 season, which raises the water in the southern hemisphere every year. The MY28 GDS, coincident
33 with the southern summer solstice, demonstrated the most massive increase of water abundances
34 at high altitudes through the observed years.

35

36 **Key points:**

37 1) Through eight Martian years observed by SPICAM H₂O regularly reaches 40-60 km in the
38 aphelion season and 70-90 km in the perihelion season.

39 2) In southern summer, large H₂O vmrs (~100 ppm) are repeatedly observed up to 80 km,
40 confirming a seasonal impact on the hydrogen escape rate.

41 3) Out of the two observed GDS, in MY28 and MY34, the GDS of MY28, coincident with the
42 southern summer solstice, shows a larger water increase.

43

44 **1. Introduction**

45 Despite the scarcity of water vapor in the Martian atmosphere it nonetheless plays a pivotal
46 role in the planet's climate. Major efforts have been made to understand how water vapor is
47 spatially and temporally distributed over the last few decades (Montmessin et al., 2017a). Its global
48 seasonal cycle has been studied continuously for 20 years through nadir measurements of column
49 abundance by TES/Mars Global Surveyor (Smith 2004), three experiments onboard Mars-Express
50 (MEX), SPICAM (Trokhimovskiy et al., 2015), OMEGA (Maltagliati et al., 2011a) and PFS
51 (Fouchet et al., 2007; Wolkenberg et al., 2011), and by the CRISM instrument on the Mars
52 Reconnaissance Orbiter (MRO) (Smith 2009). The global annual average abundance of water
53 fluctuates around 10 pr. μm (~200 ppm, assuming mixed water), increasing at high latitudes when

54 the polar caps sublimate in spring and summer of either hemisphere. The northern hemisphere
55 reaches its maximum in summer at 75°N, exhibiting a column abundance of roughly 50 pr. μm
56 recurrent from year to year, while the corresponding southern summer maximum is half as strong
57 shows some interannual variability around 25 pr. μm .

58 In contrast to column abundance, the climatology of water vapor vertical distribution has
59 only been scarcely documented to date. The vertical distribution gauges multiple processes
60 controlling the Martian water cycle, including the condensation/sublimation leading to cloud
61 formation/collapse, the associated sedimentation and scavenging of dust particles upon which ice
62 particles form, but also the sublimation and condensation from the polar caps, the
63 photodissociation and escape processes, and surface-atmosphere exchange (Montmessin et al.,
64 2017a). The vertical distribution also appears to be a insightful tool for revealing the details of the
65 Martian circulation, in particular, the massive upwards motions that permit water to access
66 altitudes higher than 100 km in some extreme cases (Shaposhnikov et al., 2019; Neary et al.; 2020).

67 By analogy with the Earth's troposphere, the vertical distribution of water on Mars is
68 assumed to be controlled by vapor pressure, which is solely determined by temperature. The
69 altitude level at which water vapor should condense in theory, the hygropause, rises from 10–20
70 km in the [40°S,80°N] latitude range during the colder aphelion period to 30–60 km in the
71 [80°S,40°N] latitude range during the warmer perihelion period (Richardson and Wilson, 2002a;
72 Montmessin et al., 2004). . Saturation and further condensation of water vapor at 10 km in northern
73 summer is responsible for the formation of the Aphelion cloud belt (ACB), which regulates the
74 transfer of water between hemispheres. Condensation blocks water in the northern tropics where
75 the seasonal low-level air mass convergence turns into the rising branch of the Hadley cell (Clancy
76 et al., 1996; Richardson and Wilson, 2002a; Montmessin et al., 2004). The seasonal trend of
77 saturation altitude can be inferred from column-integrated H₂O measurements assuming a known
78 temperature distribution, see, e.g., Trokhimovskiy et al. (2015).

79 However, the concept of saturation-controlled water holding capacity is presently undermined
80 by several lines of evidence. First, SPICAM solar occultation profiles measured in Martian Year
81 (MY) 29 at $L_s=60\text{--}110^\circ$ show considerable amounts of H_2O above the hygropause, at 30-40 km
82 (Maltagliati et al., 2011b). This water vapor, therefore, reached a supersaturated state with a
83 pressure several times (up to 10) higher than the saturation pressure. Clancy et al. (2017) have also
84 found strong, though indirect, hints of the persistence of supersaturation in water at $L_s=60\text{--}140^\circ$
85 using as a proxy the oxygen singlet emission measured by CRISM. The full second half of MY34
86 ($L_s=165\text{--}360^\circ$) observed with ACS (Atmospheric Chemical Suite) onboard the Trace Gas Orbiter
87 (TGO) shows that water vapor supersaturation is deep and ubiquitous (Fedorova et al., 2020a),
88 implying that water can be transported to the upper atmosphere much more easily than previously
89 assumed. High degrees of supersaturation were also predicted in simulations of water ice cloud
90 microphysics inside a Global Climate Model (Navarro et al., 2014).

91 Water vapor was observed at altitudes of up to 80-100 km during the Global Dust Storms
92 (GDS) of 2007 (MY 28) and 2018 (MY34) with a volume mixing ratio in the 50-100 ppmv range
93 (Fedorova et al., 2018, 2020a; Heavens et al., 2018; Vandaele et al., 2019; Aoki et al., 2019). The
94 region above 60–80 km is essential in the fate of water on Mars, as the thinner CO_2 atmosphere no
95 longer screens the solar UV light. Water photolysis here thus provides the main source of hydrogen
96 atoms in the upper atmosphere, enabling their subsequent escape to space (Chaffin et al., 2017).

97 Variations in the Martian hydrogen escape rate by an order-of-magnitude were inferred from
98 Lyman- α airglow observations during the MY28 GDS by Mars Express and the Hubble Space
99 Telescope (Chaffin et al., 2014; Clarke et al., 2014). Photochemical modeling tied the enhanced
100 escape rates with the presence of water molecules at 60-80 km during the GDS (Chaffin et al.,
101 2017; Krasnopolsky et al., 2019). More hydrogen corona observations suggested that not only dust
102 storms but the perihelion season as a whole, even without a major dust event, might cause the
103 escape rate variation (Bhattacharya et al., 2015). Instead, during perihelion and outside a GDS,
104 significant amounts of water vapor have been observed in the southern hemisphere during MY29

105 (Ls=240–260°) above 60 km (Maltagliati et al., 2013), and by ACS in MY34 (Ls=270–290°, in
106 between the dust events) up to 100 km (Fedorova et al., 2020a). GCM simulations explain the
107 seasonal upward transport of water vapor by atmospheric dynamics lifting water in the upward
108 branch of the pole-to-pole meridional circulation cell. Shaposhnikov et al. 2019 demonstrated the
109 upward water flux maximizes around perihelion between Ls= 220° and 300°, in a symmetric
110 fashion about Ls=260° when the global mean temperature reaches annual maximum and the
111 circulation is the most intense because of hemispheric dichotomy (Richardson and Wilson, 2002b).
112 Even at this moment, the upward water flux is minimal at ~60 km, where water can penetrate into
113 the upper atmosphere only in the region between 20°S and 70°S and transported further upward
114 and across latitudes northward. At other seasons this flux is negligible.

115 Better seasonal coverage of the vertical water distribution in the lower atmosphere (<100 km)
116 of Mars and of the Lyman- α airglow in the upper atmosphere (>200 km) should help disentangle
117 the respective roles between the intense but rare GDS and the weaker but recurrent seasonal
118 increase of dust in modulating the escape of hydrogen from Mars. The monitoring of water vapor
119 profiles is also necessary to understand whether the high-altitude water, which constitutes a small
120 fraction of the whole atmospheric water, is an essential component in the transport of water across
121 the globe.

122 Here, for the first time, we present observations of H₂O vertical distribution in the Martian
123 atmosphere obtained by SPICAM infrared spectrometer onboard Mars Express during eight Mars
124 years, from MY27 to MY34, including two GDS. The instrument measures the water density and
125 mixing ratio at the Mars limb in solar occultations (Korablev et al., 2006; Fedorova et al., 2009,
126 2018; Montmessin et al., 2017b; Maltagliati et al., 2011b, 2013). The details of the SPICAM
127 observations and the data analysis are described in sections 2 and 3, respectively. The retrieved
128 climatology is presented in section 4, where we discuss the average dataset with a focus on the
129 second, dusty half of the Martian year, from Ls=180° to 360°, and the two GDS.

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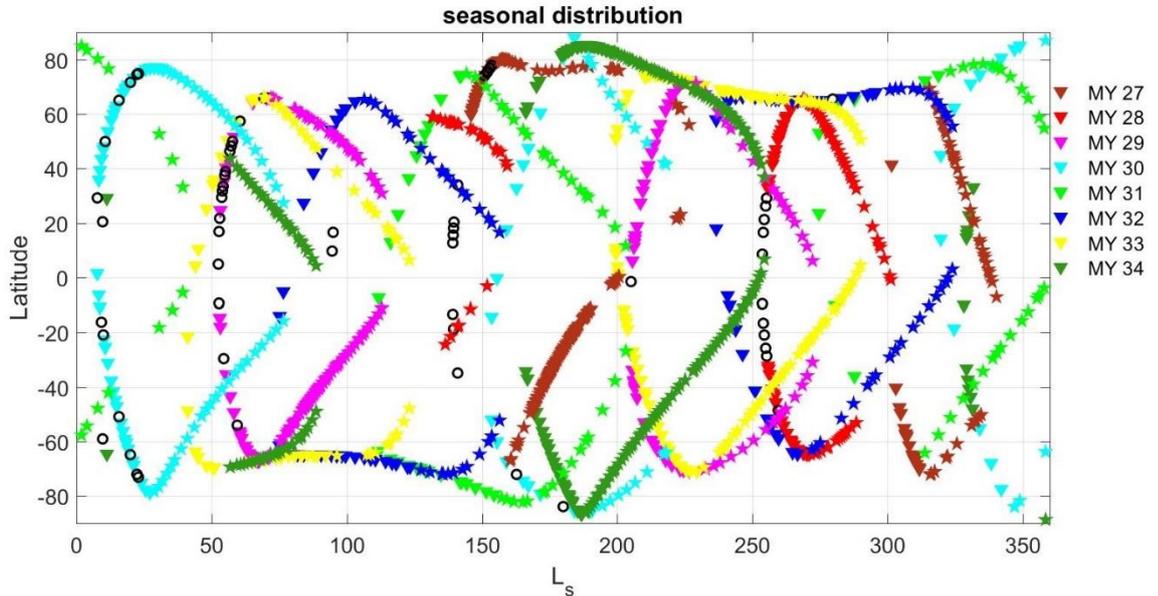
2. SPICAM observations.

SPICAM IR is an AOTF (Acousto-Optic Tunable Filter) infrared spectrometer. It works in the spectral range of 1–1.7 μm with a spectral resolution of 3.5–4 cm^{-1} . Near the H_2O 1.37- μm band the resolving power corresponds to ~ 2000 . The spectrometer operates both in nadir and solar occultation modes; the present work is based on occultation data. Two detectors of the spectrometer working in orthogonal polarizations have different signal-to-noise ratios (SNR). This corresponds to ~ 120 in channel 0 and ~ 220 in channel 1 for a pure solar spectrum outside the atmosphere in the range of 1.2–1.55 μm . Thus we use the data of channel 1 for the analysis. With the FOV of 4.2 arcmins (0.07°) the vertical resolution in occultations varies from 1 to 12 km (4–5 km on average) due to the elliptical orbit of the Mars-Express spacecraft. More details of the instrument and its calibrations related to occultations can be found in papers by Korablev *et al.* (2006a, b) and Fedorova *et al.* (2009).

SPICAM can measure the H_2O and CO_2 density and the aerosol density and properties. The AOTF implies a sequential acquisition of spectra, but the measured spectral range can be arbitrarily chosen. In solar occultation we use 609 spectral points in the range of 1.34–1.47 μm to record the strongest CO_2 1.43- μm band, the H_2O 1.37- μm band and 55 points distributed between 1 and 1.7 μm to measure transmission outside gaseous absorption bands (“reference wavelengths”) and characterize the optical properties of aerosols (see also Fedorova *et al.*, 2009, 2018). One spectrum is recorded in 4 s. Taking into account that the vertical projection of the spacecraft speed on the limb varies from 0.5 to 3 km/s, the change in altitude over one spectrum varies from 2 to 10 km depending on occultation.

From 2005 to March of 2019 Mars Express has completed 18 occultation campaigns. Figure 1 shows the coverage of solar occultations for eight Martian years from MY27 to MY35. This dataset includes ~ 1500 occultations. During the two first occultation campaigns in MY27 the spectral coverage was not optimized, which meant that only the H_2O density could be retrieved, but not the CO_2 density. Through the whole dataset, $\sim 15\%$ of data are not suitable for the retrieval

157 of gaseous density profiles due to pointing instabilities, problems with the reference solar
 158 spectrum, etc.



159
 160 Figure 1. Seasonal-latitudinal coverage of SPICAM observations in solar occultations for eight
 161 Martian years. The open circles mark faulty observations inapplicable for retrieval. Morning
 162 observations are marked by triangles and evening observations by stars.

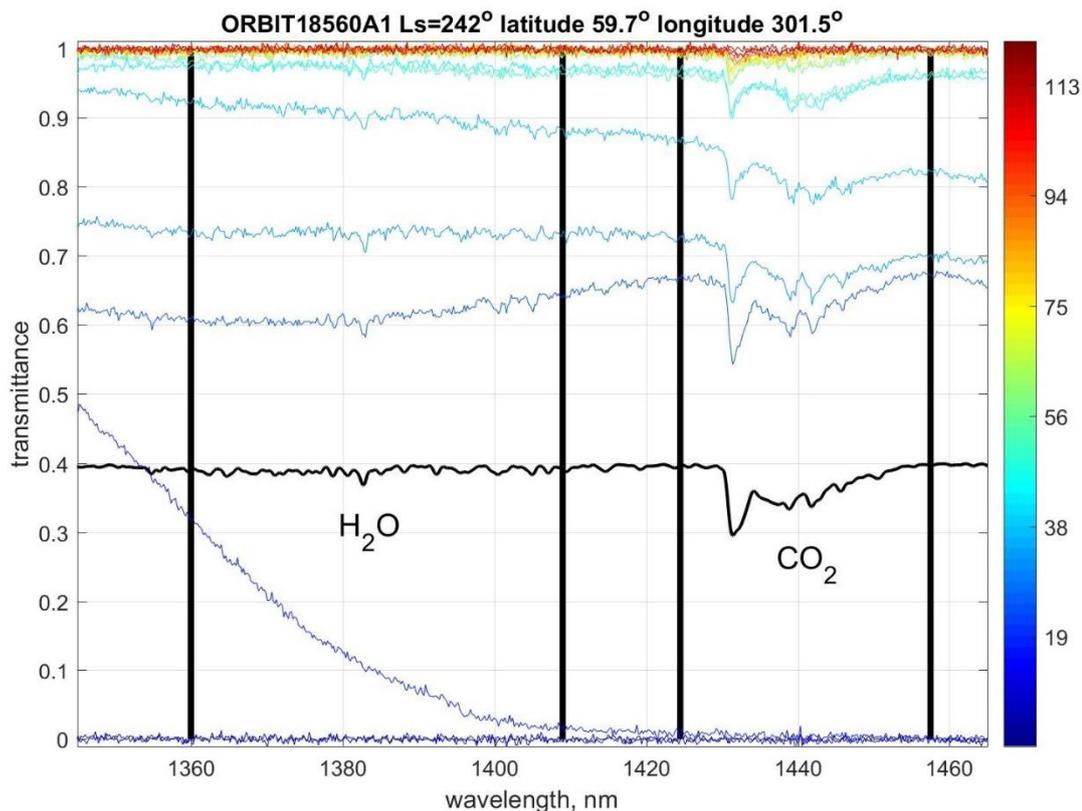
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164 3. Data processing

165 The details of the SPICAM IR data processing in occultations is described in Fedorova et al.
 166 (2009, 2018) and Maltagliati et al. (2011, 2013). A summary of the retrieval process and a few
 167 recent amendments to it are presented below.

168 Occultation observations are self-calibrated, yielding transmittance of the atmosphere from the
 169 ratio of a spectrum through the atmosphere to the solar reference collected high enough where the
 170 atmosphere negligibly absorbs. Above 120 km for any season and location the CO₂ 1.43 μm
 171 absorption band is too weak to be detected by SPICAM; we therefore averaged the reference
 172 spectrum between 120 and 170 km. Fig. 2 shows an example of transmittances obtained in orbit
 173 18560 of MY34.

174



175

176 Figure 2. SPICAM transmittance spectra measured between 0 and 120 km for orbit 18560A1
 177 (Ls=242°, 59.7°N, 301.5°E) in the spectral range of 1346–1462 nm. The color indicates the
 178 average altitude of a spectrum. Data of detector 1 are shown. The black curve is the model
 179 spectrum with the CO₂ and H₂O absorption bands at the target altitude of 20 km. Solid lines bound
 180 the spectral ranges used for the retrievals. A broad and shallow absorption appearing at 1350–
 181 1390 nm (~0.6 transmission level) is likely a signature of a thin detached aerosol layer, manifested
 182 as a change of the slant optical depth during the spectrum record.

183

184 Because each spectral point in an occultation spectrum is measured sequentially, each point
 185 corresponds to a different altitude. The aerosol opacity, which changes with altitude, results in an
 186 artificial spectral slope (see Fig. 2). To correct it, we normalized the spectra around the studied
 187 absorption bands to the continuum outside the bands. For CO₂, the continuum is computed with a
 188 linear interpolation between averages of two intervals on both sides of the band, 1423–1426 nm
 189 and 1456–1459 nm. The H₂O band at the short-wavelength boundary extends beyond the acquired

190 spectral window, and the continuum is estimated in between the absorption lines within the band
191 as described in Maltagliati et al. (2013).

192 To reduce the noise in measured spectra, we applied the Savitzky-Golay smoothing filter
193 (Savitsky and Golay, 1964) to ensure a strong noise reduction while preserving the spectral
194 resolution. No averaging of spectra along the occultation sequence is performed, because of the
195 too coarse vertical sampling, with typically 5–25 useful spectra during an occultation between 0
196 and 100 km. We also limit the H₂O fit to the strongest part of the band, between 1360 and 1409
197 nm.

198 The whole profile of the local H₂O or CO₂ density is retrieved using the Levenberg-Marquardt
199 iterative algorithm (Levenberg, 1944; Marquardt, 1963). The Tikhonov regularization is applied
200 after the fit, in order to smooth the profile and minimize the errors (Ceccherini, 2005; Ceccherini
201 et al., 2007). The uncertainty in the densities is given by the covariance matrix of the solution
202 errors. The H₂O detection limit is estimated as $7\text{--}9\times 10^9$ molecules/cm³. This detection limit
203 corresponds to the volume mixing ratio (vmr) accuracy better than 1 ppm below 35 km, better than
204 10 ppm at 50–55 km and better than 70 ppm at 80 km. The measurement accuracy varies with
205 season and location due to variations of atmospheric density. The CO₂ detection limit is below
206 10^{12} molecules/cm³ which corresponds to the pressure level of $2\times 10^{-6}\text{--}6\times 10^{-5}$ mbar. The 1.43 μm
207 absorption band is visible up to 110–115 km depending on location and season.

208 A line-by-line synthetic spectrum of gaseous absorption was calculated using the HITRAN
209 2012 spectroscopic database (Rothman et al., 2013). There are only a few laboratory measurements
210 of the CO₂-broadened H₂O half-widths in this band for some lines (Langlois et al., 1994). Based
211 on theoretical calculations by Gamache et al., (1995) and measurements in thermal IR by Brown
212 et al., (2007), we correct the air-broadening half-widths for the CO₂-broadening in the Martian
213 atmosphere by multiplying by a factor of 1.7 (see also discussion in Fedorova et al., 2010). In case
214 of occultation measurements, the Doppler broadening begins to dominate above 30–40 km and the
215 sensitivity to this coefficient is minimal.

216 Information about the pressure and temperature profile in the Martian atmosphere is a critical
217 aspect of SPICAM's water retrieval, to simulate spectra for the forward model, as well as to
218 compute the mixing ratio. As described in Fedorova et al. (2018), while the retrieved H₂O density
219 is not very sensitive to the temperature profile, the CO₂ density could be mistaken by a factor of 2
220 when the temperature deviates by as much as 50 K. This translates into a 50% error in the water
221 mixing ratio. For the forward model we used the temperature and pressure profiles from the recent
222 European Martian Climate Database (EMCD) [<http://www-mars.lmd.jussieu.fr/>, version MCD
223 V5.3; Forget *et al.* 1999; Millour *et al.* 2019]. To mitigate the differences between the model and
224 the observations we choose the individual Mars years scenarios available in the MCD for MY 24
225 through 34, which use the observed daily dust column maps for each year to constrain the dust
226 distributions in the model and hence reproduce observations to the greatest extent possible
227 (Montabone et al., 2015, 2020). These scenarios reproduce the temperature profiles by year
228 reasonably well, with inaccuracy being far below 50K outside of a GDS. In the case of CO₂, the
229 density profile predicted by MCD was used as an *a priori* vector in the minimization procedure.
230 For H₂O, we used a constant volume mixing ratio of 1 ppmv at all altitudes as an initial assumption.
231 Examples of H₂O fits can be found in Fedorova et al. (2009; 2018).

232 Even if individual dust scenarios of MY25, 28, and 34 in the MCD include the extreme
233 global dust events, modeling of a GDS remains challenging. As shown in Fedorova et al. (2018),
234 during the MY28 GDS a difference between MCD temperature profiles and measured MCS/MRO
235 profiles sometimes reached 40 K. Such a bias can be critical for our retrievals. We resolved to use
236 MCS temperature profiles (Kleinböhl et al., 2009) close in time/space for these special cases when
237 available. MCS measures at 03:00 and 15:00 Martian local time, out of phase with the Mars-
238 Express occultations, but many observations within 1° of Ls and 2° of latitude and 5° of longitude
239 could be found. The selection of the MCS data was done at local times closest to the SPICAM
240 observations and performed additional checking using the daily temperature cycle of the MCD to
241 find the difference between morning-evening SPICAM observations and day-night MCS profiles,

242 as detailed in Fedorova et al., 2018. The share of available collocated MCS temperature profiles
243 during two GDS (in MYs 28 and 34) used to retrieve the H₂O densities is high, reaching 90% of
244 all observations.

245 The vertical profile of the H₂O volume mixing ratio (f_{H_2O}) is the ratio between the water
246 vapor number density and the atmospheric number density. To find the atmospheric density profile
247 we divide the CO₂ profile by the CO₂ mixing ratio, assumed to be 0.9543. As the CO₂ and the H₂O
248 bands were not acquired at the same altitude (as described above) we find the atmospheric density
249 at the H₂O altitude with a log-linear interpolation. The error in the H₂O mixing ratio is given by
250 the quadratic sum of the uncertainties of the two density profiles. The main systematic error on
251 H₂O vmr results from the temperature profile uncertainty, in particular during the dusty part of the
252 Martian year, due to a strong sensitivity of the CO₂ band, as described above. The efforts taken to
253 minimize this uncertainty reduce the error in the H₂O vmr to <20–30%. All the details of the
254 retrieval and the associated sensitivity analysis are presented in previous works (Maltagliati et al.,
255 2011b, 2013, Fedorova et al., 2009, 2018).

256

257 **4. Results**

258 **4.1 H₂O density distribution**

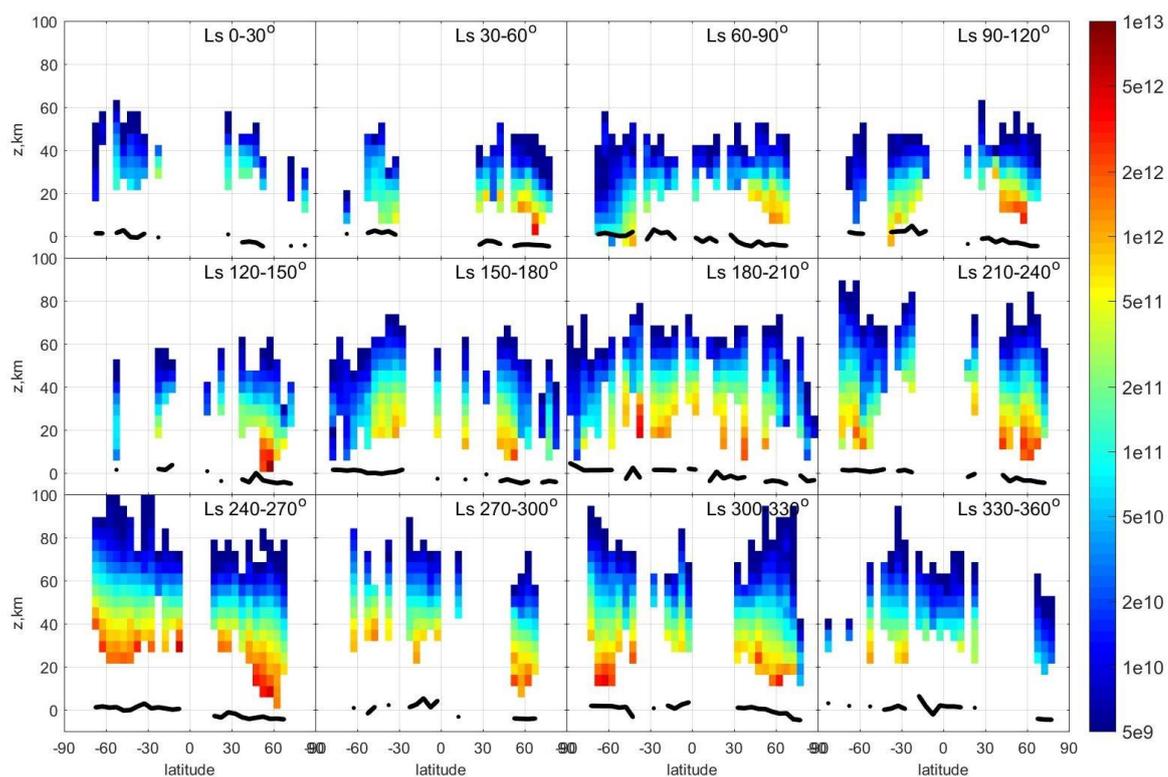
259 The SPICAM water density profiles have been obtained for eight Martian years from MY 27
260 to 34 with two occultation campaigns per year. The vertical distribution of water vapor is highly
261 variable with season and altitude. To obtain the latitude vs. altitude distribution seasonal
262 distribution through regular years without dust storms, we averaged all observations into 12×30°
263 Ls bins (Figure 3) and then averaged each hemisphere separately to compare their respective Ls
264 vs. altitude distributions (Figure 4). The two GDS of MY28 (Ls=267–302°) and MY34 (Ls=190–
265 230°) characterized by a high elevation of water vapor are excluded from the averaged maps.

266 The detected density varies from $5 \times 10^9 \text{ cm}^{-3}$ to 10^{13} cm^{-3} . Water vapor is less abundant and
267 more confined into the lower atmosphere during northern spring/summer, as expected, due to the

268 colder aphelion climate and correspondingly the lower level of the hygropause. Noticeable
269 amounts of water are observed at 30-40 km, while profiles can be traced as high as 40–60 km in
270 the aphelion season at $L_s=30\text{--}120^\circ$. Around perihelion, at $L_s=240\text{--}270^\circ$ the highest water is
271 observed at 70–90 km. To first order then, water transport to higher altitude is dictated by
272 atmospheric temperature as the Aphelion-to-Perihelion transition corresponds to $>20\text{K}$
273 temperature change (Smith et al., 2017).

274 At $L_s=30\text{--}90^\circ$ during the northern spring-summer, the maximal water density was observed in
275 the middle and high northern latitudes corresponding to the sublimation at the northern polar cap
276 and subsequent water release. In the southern equinox ($L_s=150\text{--}210^\circ$) timeframe, the water
277 distribution is very symmetric between the two hemispheres with low values at high latitudes.
278 Sparse detections at $L_s=0\text{--}150^\circ$ are explained by the confinement of the water vapor near the
279 surface, associated with a less confined aerosol layer and Aphelion intertropical clouds that both
280 impair detection.

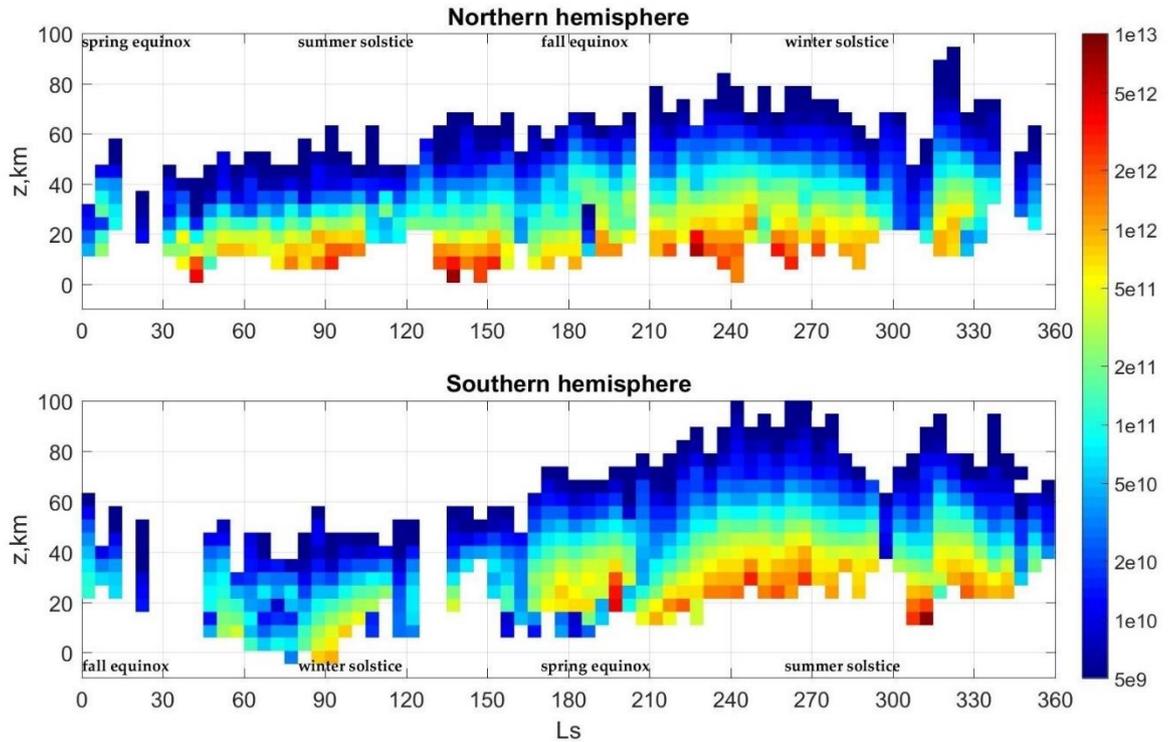
281 Figure 5 shows the maximal altitude where the H_2O density reached 10^{10} (close to SPICAM's
282 detection limit, but still dependably measured). The black open circles mark orbits where no water
283 was detected at all, meaning the density at all altitudes was below $5\text{--}7\times 10^9\text{ cm}^{-3}$. This happened
284 for some of the low-latitude orbits associated with the aphelion clouds at $L_s=0\text{--}150^\circ$. Detections
285 are also absent at polar latitudes, poleward of 60°S in southern winter ($L_s=20\text{--}180^\circ$), and poleward
286 of 60°N in northern spring ($L_s=330\text{--}30^\circ$) and autumn ($140\text{--}200^\circ$). The 10^{10} cm^{-3} detection
287 threshold altitude varies from 15 to 40 km (30 km on average) in the aphelion season, rising to 60-
288 90 km in the perihelion season.



289

290 Figure 3. The H₂O density distribution (in cm⁻³) with altitude and latitude for 12 Ls bins
 291 averaged over the MY27-34 timeframe. The altitudes and latitudes have been arranged into 5km
 292 × 5° bins. The dataset for the GDS of MY28 and 34 were excluded. The y-axis shows the altitude
 293 above the areoid. Black curves mark averaged MOLA altimetry for the bins.

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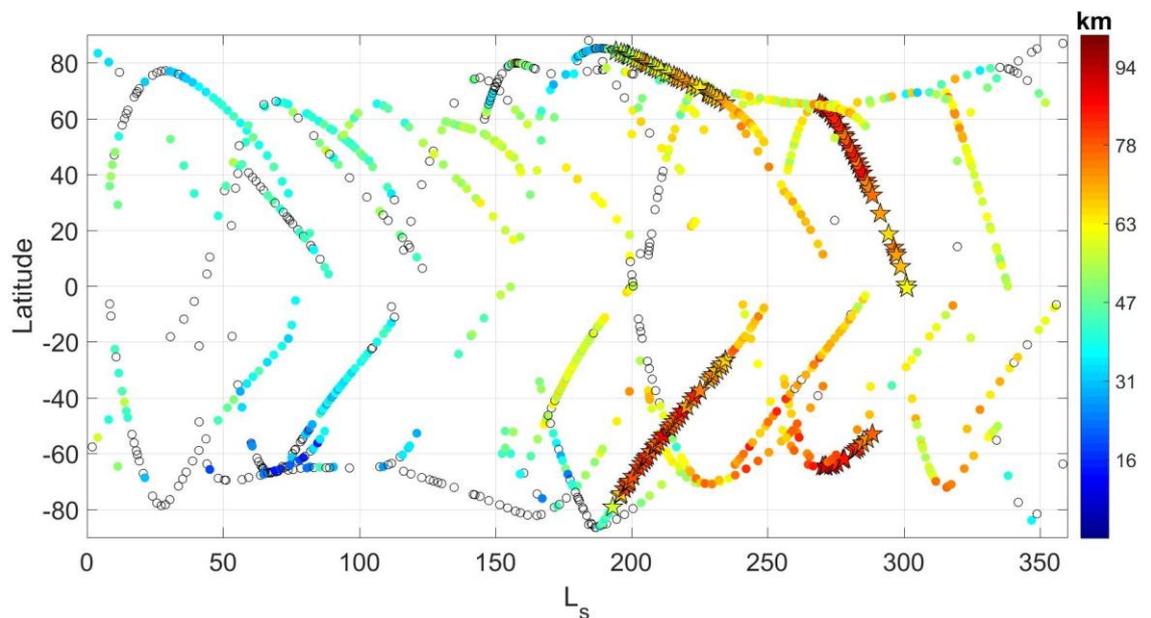
296 Figure 4. The H₂O density distribution (in cm⁻³) with altitude as a function of Ls for the northern
 297 and southern hemispheres. The Ls scales are shifted to align summer and winter seasons in both
 298 hemispheres. Altitudes and Ls are arranged into 5 km × 5° bins. Two GDS of MY28 and 34 are
 299 excluded from the binning.

300 In contrast, the warmer southern spring appears associated with higher vertical mobility of
 301 water vapor. From Ls=210° the maximum of water density moves to middle-to-high southern
 302 latitudes, marking the dissipation of the polar cap. The water amount increases after Ls = 240° to
 303 solstice, when the density reaches 10¹² cm⁻³ above 40 km and water extends to 90 km (Fig. 3 and
 304 4). Several periods when the high-altitude water is observed in the dusty season can be identified
 305 in Figure 5. Two periods, marked by stars, are related to the GDS of MY28 and 34 at Ls=190-240°
 306 and 270-300°, respectively. During these times, water attains altitudes up to 90 km in both
 307 hemispheres. A pronounced solstitial maximum at Ls=240–300° corresponds to water lofted up
 308 to 70-90 km in the southern hemisphere every year. In the northern hemisphere, water reaches only

309 60-80 km and thus the picture looks asymmetric around the equator in the corresponding panel of
 310 Fig. 3.

311 We also note a moderate, compared to perihelion and GDS, increase of water up to 60–80 km
 312 at $L_s = 315\text{--}330^\circ$ in both hemispheres (Fig. 4), at least for three occultation campaigns MY27, 31
 313 and 32 (Fig. 5). This increase is correlated with regional dust activity, which occurs at this L_s
 314 almost every year (Kass et al., 2016). This regional dust activity started between $L_s = 305^\circ$ and 320°
 315 and lasted $3\text{--}15^\circ$ of L_s . Interestingly, water increased in both hemispheres in middle to high
 316 latitudes from 40° to 70° in accord with the recurrent northern temperature response (Kass et al.,
 317 2016). A similar increase of water was observed during the MY34 GDS by TGO (Aoki et al.,
 318 2019; Fedorova et al., 2020a).

319



320
 321 Figure 5. The altitude of water density of 10^{10} cm^{-3} (close to SPICAM's detection limit). Black
 322 open circles mark orbits where water is below the detection limit everywhere in the profile. Stars
 323 correspond to the GDS of MY28 and 34.

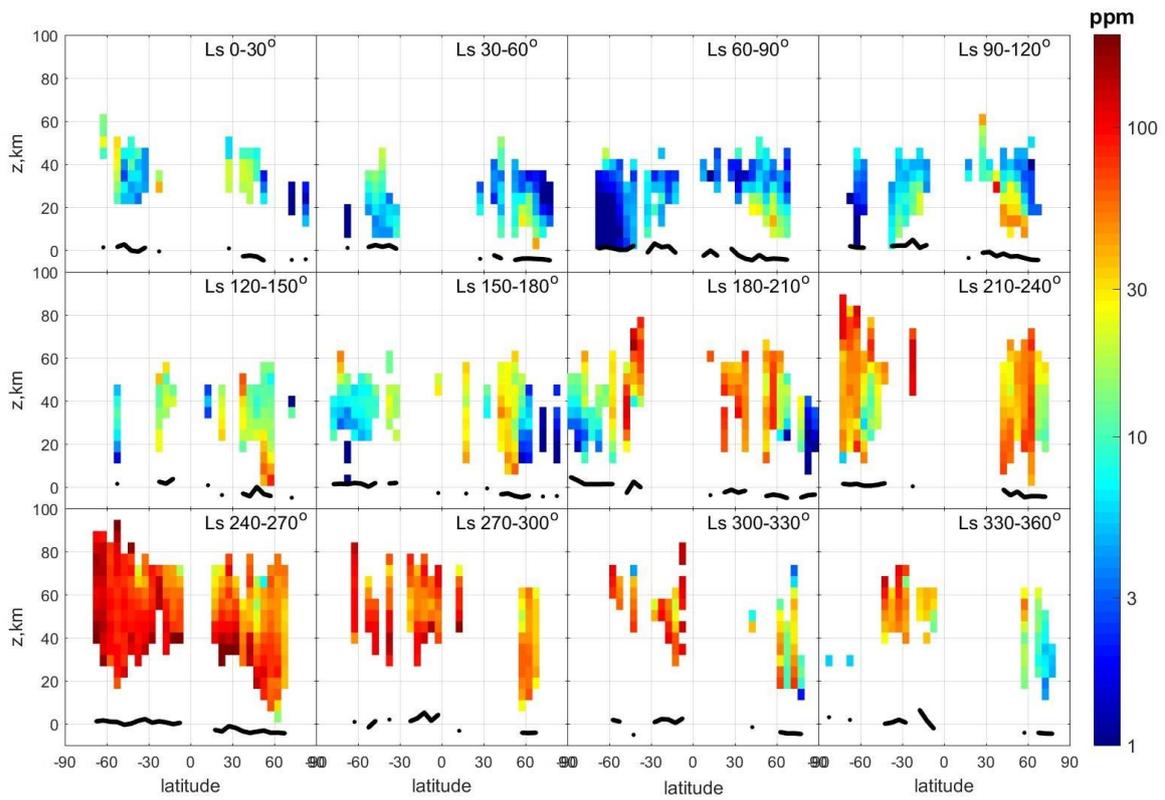
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325 **4.2 Mixing ratio**

326 Figures 6 and 7 show the H₂O vmr latitude-altitude cross-sections for 12 seasonal bins and for
327 the two hemispheres, respectively, similar to Figures 3 and 4. In the aphelion season at L_s=30–
328 120°, the H₂O vmr does not exceed 40 ppmv. A minimum of water is found in middle-to-high
329 southern latitudes during southern winter (L_s=60–90°), where the vmr was below 1–2 ppmv. The
330 confinement of water vapor towards the surface during the aphelion season suggested by the
331 present observations (as the lowest part of the profile is inaccessible to occultation) appears in
332 agreement with the conclusions drawn from nadir observations (Smith, 2004; Trokhimovsky et
333 al., 2015). Nonetheless, a sharp decrease of the mixing ratio from 50 to 5 ppmv occurs in the
334 middle northern latitudes at L_s=30–120° (the most prominent at L_s=90–120°), indicating a rise of
335 the hygropause from 15 km at 60°N to 25–30 km at 40°N.

336 The water mixing ratio in low-to-middle latitudes starts to grow in the middle atmosphere of
337 both hemispheres after L_s=180°, reaching 50 ppmv between 20 and 60 km. Higher northern and
338 southern latitudes are exempt from this moistening and remain dry with vmr ≤ 5–10 ppmv. In the
339 L_s=210–240° period, humid air progressively fills all latitudes, and reaches altitudes up to 70–80
340 km. The most abundant water (>100 ppm) can be observed at L_s=240–270° in both hemispheres,
341 spanning the whole middle atmosphere (from 30 to 70–80 km) exhibiting an uniform distribution.
342 At high southern latitudes the water layer expands vertically higher than 90 km. A similar pattern
343 is observed at L_s=270–300°, even if our occultation dataset does not cover this season well enough.
344 The mixing ratio decreases only after L_s~300° in the north, as tracked by the altitude of the 10¹⁰
345 density level in Figure 5. However, the mixing ratio in the low-to-mid southern latitudes still
346 reaches 50–100 ppmv at 70 km. Remains of the southern summer water maximum can still be
347 spotted at low latitudes throughout the L_s=330–360° period.

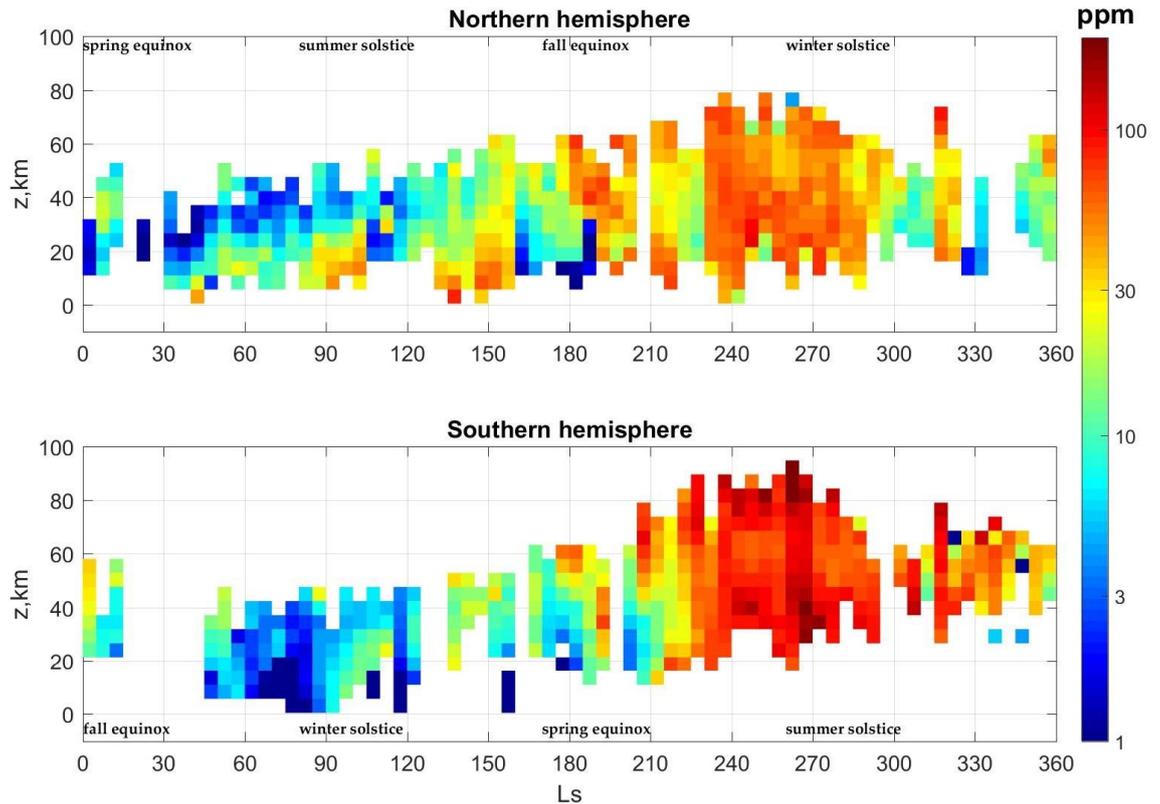
348



349

350 Figure 6. Same as Figure 3 except the H₂O mixing ratios are displayed instead of the number
 351 density.

352



353

354 Figure 7. Same as Figure 4 except the H₂O mixing ratios are displayed instead of the number
 355 density.

356 *4.3 Comparing years with and without GDS during the dusty season*

357 The SPICAM survey period includes the two GDS that happened in 2007 (MY28) and
 358 2018 (MY34). Wang and Richardson (2015) described the development of the MY 28 GDS as
 359 observed by MARCI (MARs Color Imager) on MRO. The storm reached its global phase on
 360 $L_s=269.3^\circ$ (13 sols after onset) at southern middle-to-high latitudes between $35\text{--}80^\circ\text{S}$ (Cantor et
 361 al., 2008). By the end of the expansion phase ($L_s\sim 275.2^\circ$), the dust encircled the planet from
 362 approximately 90°S to 58°N , reaching an optical depth of 4.4–4.6 (Smith, 2009b; Lemmon et al.,
 363 2015). The decay phase started at $L_s\sim 310^\circ$ and ended at $L_s\sim 310\text{--}320^\circ$.

364 The 2018 GDS has been reviewed by Guzewich et al. (2019), Montabone et al. (2020), and
 365 references therein. The storm started regionally in the northern hemisphere in the second half of
 366 May 2018 ($L_s\sim 186^\circ$) and then merged by the beginning of June ($L_s\sim 190^\circ$) with another regional

367 dust lifting occurring independently in the southern hemisphere. It became global by the middle
368 of June ($L_s \sim 197^\circ$). This most active period of the GDS lasted until the beginning of July (L_s
369 $\sim 205^\circ$) followed by a long decay, which ended in the middle of September ($L_s \sim 240\text{--}250^\circ$).

370 The process of dust lifting was uneven across the globe for both GDSs, and even at the peak
371 of the storm northern latitudes ($>50^\circ\text{N}$) and southern polar areas were mostly clear of dust and
372 suitable for solar occultations (Montabone et al., 2015; 2020). The intercomparison of the two
373 GDS (Smith, 2019; Wolkenberg et al., 2020) showed that even though they happened in different
374 seasons, their globally-averaged peak optical depths were similar. The MY 28 GDS reached dust
375 optical depth greater than 1.5 (at $9\ \mu\text{m}$) rapidly, within $10^\circ\text{--}15^\circ$ of L_s , and then maintained this
376 dust loading for an extended period, before decay. According to THEMIS observations the MY
377 34 GDS was characterized by a more gradual build-up with a two-stage increase in dust optical
378 depth. Dust optical depth increased rapidly at the start of the storm at the same rate as during MY
379 28 but then remained nearly constant for several degrees of L_s before increasing again to its
380 maximum value (Smith, 2019).

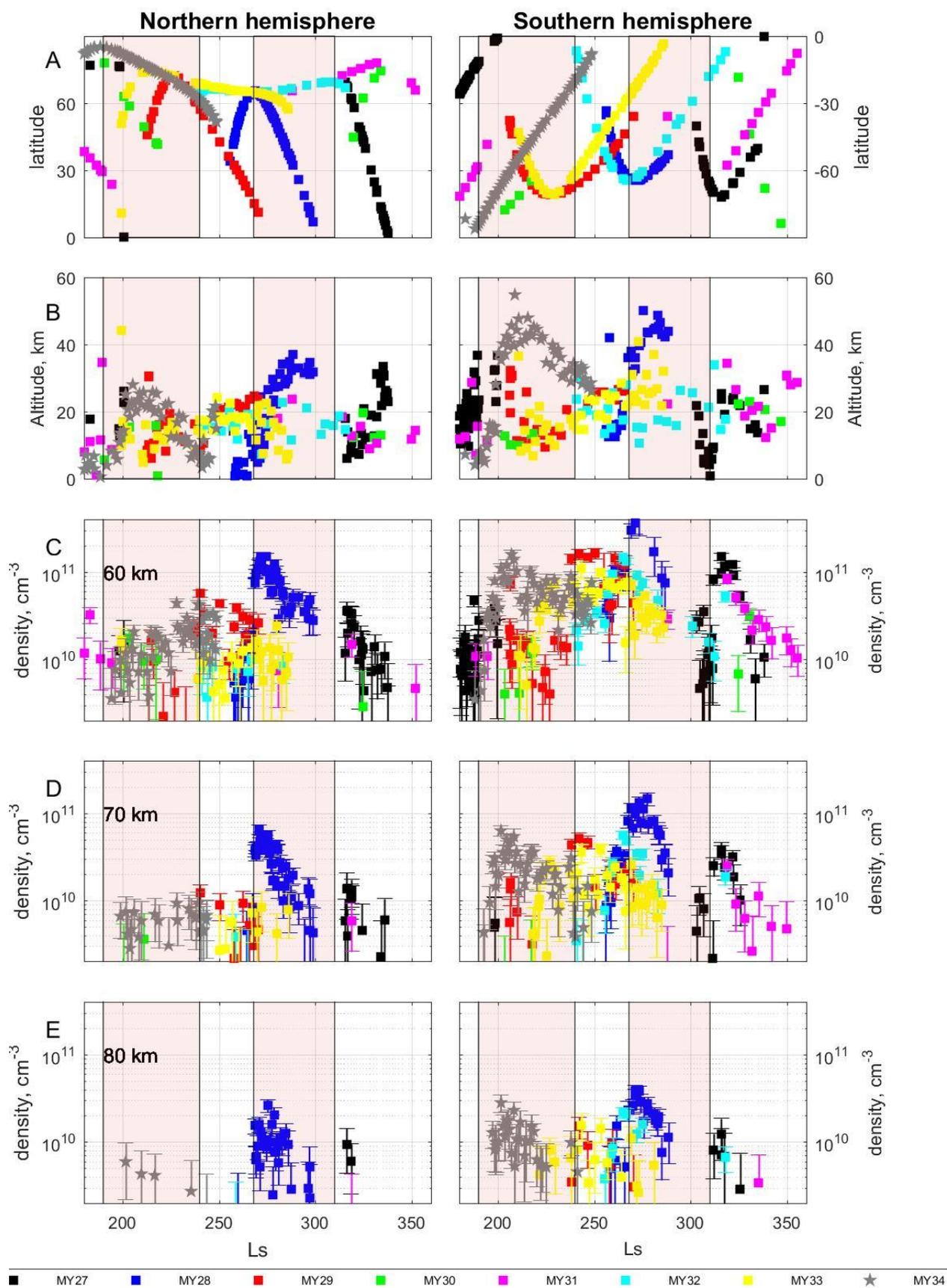
381 The water vapor distribution during MY28 was described in detail by Fedorova et al.
382 (2018). Here we compare the behavior of water during the two GDS, highlighting the contrast with
383 regular perihelion seasons, unaffected by GDS. Figures 8 and 9 present the H_2O density and mixing
384 ratio distribution for altitudes of 60, 70 and 80 km and the lowest usable tangent height for
385 SPICAM occultations, limited by transmission levels below this point being <10 percent which
386 corresponds to aerosol optical depth of ~ 3 on the line of sight and indicating the aerosol loading.
387 In the northern hemisphere from $L_s=268^\circ$ to $L_s=285^\circ$ the H_2O density increased by an order of
388 magnitude at 60-80 km (fig.8). During the dust storm, the profiles extended up to 80 km, with a
389 H_2O mixing ratio ≥ 100 ppmv (fig.9). We note two maxima of the H_2O density. The first and largest
390 maximum was observed above 60°N , from $L_s=269$ to 275° and does not directly correlate with
391 the aerosol loading. It likely relates to the downwelling branch of the meridional circulation,
392 transporting water from the southern to the northern hemisphere, which intensified during the

393 GDS. The second and smaller maximum at $L_s=280^\circ$ coincides with the high dust loading at middle
394 northern latitudes. In the South the water density at 50–80 km increased by a factor of 4–5 with a
395 mixing ratio >100 ppm, again correlated with the aerosol extension.

396 During the MY34 GDS, SPICAM observations in the northern hemisphere began at very high
397 latitudes, $>75^\circ\text{N}$. Only from $L_s=200^\circ$ was a water density of $>10^{10}$ cm^{-3} detected with a mixing
398 ratio up to 100 ppmv at altitudes of 60-70 km and well-correlated with increase of aerosol loading
399 (Fig. 8). Compared to the MY28 GDS, water density at 60 to 80 km was one order of magnitude
400 lower. In the southern hemisphere, water density peaked at 10^{11} cm^{-3} at 70 km, and 5×10^{10} cm^{-3}
401 at 60 km, with mixing ratios of 50-100 ppmv but still 2–3 times lower compared to MY28.

402

403



404

405 Figure 8. Seasonal distribution of observations for two hemispheres, (A), the lowest altitude

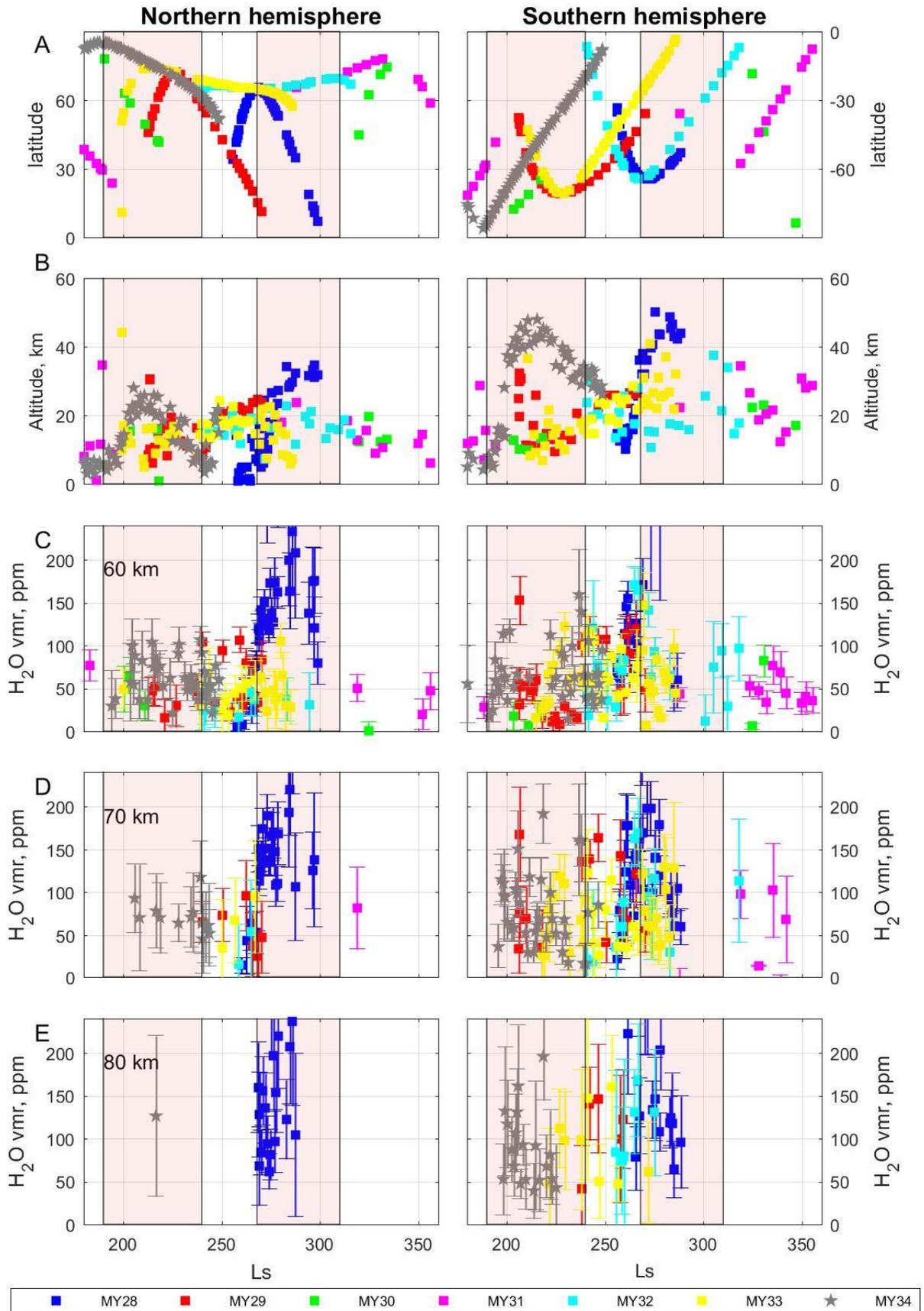
406 attained in occultation (a measure of aerosol loading), (B), the H₂O density at 60, 70 and 80 km
407 (C, D, E, respectively) for MY27-MY34. The altitudes above the areoid are presented. The two
408 GDS are marked by the pink shaded Ls periods.

409

410 Compared to a “regular” year, H₂O at 60-70 km during MY28 is several times more abundant
411 in the North. In the South, high H₂O density is also observed in MY29, 32, and 33, but still only
412 half as large as in MY28. The density increase for the same Ls during MY34 in the southern
413 hemisphere exceeds other years by 2-3 times, comparable with perihelion regular density, but
414 lower than during the MY28 GDS. Despite the observational selection, we conclude that the MY28
415 GDS was exceptional in terms of delivering water to high altitudes (80 km), exceeding at least by
416 2 times the perihelion water maximum and the water supply during GDS of MY34.

417 As one of possible explanations is the existence of a seasonal water “pump” mechanism
418 responsible for the upward transport of water vapor suggested by Shaposhnikov et al. (2019). A
419 combination of the mean vertical flux with variations induced by solar tides facilitates penetration
420 of water through the condensation at ~60 km. This mechanism becomes efficient at high southern
421 latitudes (>60°S) at perihelion. That time the upward branch of the meridional circulation is
422 particularly strong. The recent supersaturation observation in this region (Fedorova et al., 2020a)
423 even facilitates this process. Water is lifted up by the upward branch and the meridional circulation
424 then transports it to the northern hemisphere. The GDS MY28, overlapping the upward branch of
425 the meridional circulation that happened southern summer Ls=250-280° in perihelion, made the
426 lifting of water to extremely high altitudes and transports to middle-to-high northern latitudes.

427



428

429 Figure 9. Same as in Figure 8 for the H₂O mixing ratio.

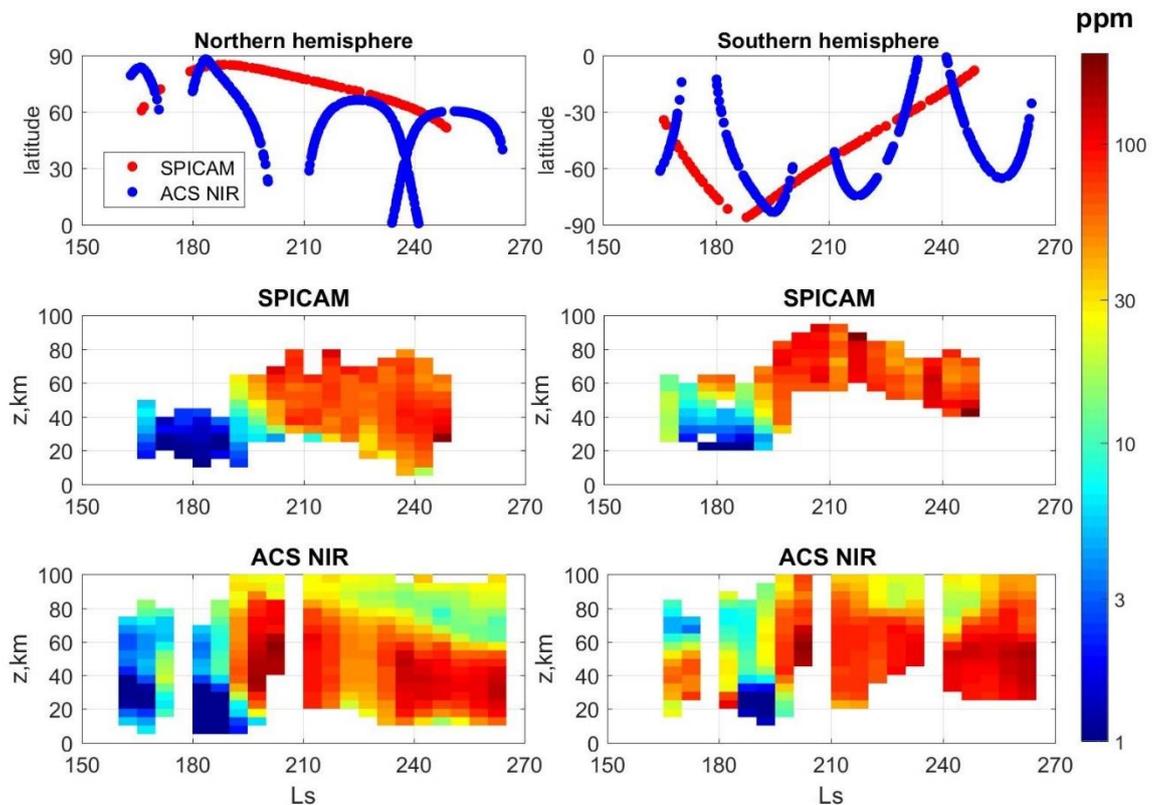
430 *4.4. The MY34 GDS as observed with SPICAM and ACS NIR onboard TGO*

431 Since April 2018, two TGO infrared spectrometers (ACS and NOMAD) perform vertical
432 profiling of water vapor. A nearly polar, not sun-synchronized circular 400-km orbit of TGO is
433 well adapted for solar occultations which began from $L_s=163^\circ$ of MY34. ACS and NOMAD
434 measure the vertical distribution of H_2O in several bands (Korablev et al., 2018; Vandaele et al.,
435 2018). Aoki et al. (2019) reported the H_2O distribution during the GDS and the regional storm at
436 $L_s=315\text{--}330^\circ$ based on NOMAD SO observations in the 2.56- μm band. Fedorova et al. (2020a)
437 reported the water and temperature for the full duration of the dusty season from $L_s=165^\circ$ to 360°
438 with detection of supersaturation of water in this period based on ACS NIR observation in the 1.38
439 μm band.

440 ACS NIR is an AOTF echelle spectrometer working in the range of 0.76–1.7 μm with spectral
441 resolving power of 25,000–27,000 (Korablev et al., 2018). In occultation, CO_2 density and
442 temperature profiles are retrieved from 1.57- μm CO_2 band and H_2O density from 1.38- μm band
443 (Fedorova et al., 2020a). Simultaneous measurement of temperature minimizes the uncertainties,
444 discussed in Section 3 for the case of SPICAM retrievals. The vertical resolution of NIR is ~ 1 km,
445 its sensitivity is better than 1 ppm between 10 and 75 km and ~ 20 ppm at 100 km, which surpasses
446 the SPICAM characteristics.

447 In order to validate the multi-year SPICAM dataset, we compare the data of SPICAM IR and
448 ACS NIR during the MY34 GDS from $L_s=160^\circ$ to $L_s=260^\circ$ in Figure 10. The data were binned in
449 the same manner for both experiments into 5 km of altitude and 5° of L_s . We note a much larger
450 altitude range, in which the ACS profiles can be retrieved. Although the occultations as observed
451 from the two spacecraft occur at different latitudes, the retrieved water trends are similar. Also,
452 the retrieved values are generally close, with the exception of cases with evident latitudinal
453 mismatch. In the northern hemisphere the increase of water to >50 ppm as observed by SPICAM
454 began at $L_s=195\text{--}200^\circ$ whereas in NIR data the growth started one bin earlier ($L_s=190\text{--}195^\circ$). This

455 shift can be explained by high latitudes sounded by SPICAM at the beginning of GDS. In the
 456 southern hemisphere, both experiments show the maximum of water at $L_s=200-220^\circ$.



457

458 Figure 10. The comparison of SPICAM and ACS NIR observations of the MY34 GDS. Seasonal
 459 distribution with altitude is shown for two hemispheres. Altitudes and L_s are arranged into
 460 $5 \text{ km} \times 5^\circ$ bins. Top panels: Distribution of occultations for both instruments; Middle panels:
 461 SPICAM observations; Bottom panels: ACS NIR observations.

462

5. Conclusions

463

464 Here we present the first multiyear survey of water vapor vertical distribution obtained
 465 from occultation measurements performed by the SPICAM spectrometer onboard the Mars
 466 Express spacecraft. The water vapor density profiles have been retrieved for 8 Martian Years
 467 MY27–MY34 and the water vapor mixing ratio for 7 Martian Years MY28–MY34, because the
 468 CO_2 density reference was not measured during the first year (MY27). Data reveal seasonal and
 469 spatial variations, and the impact of two global dust storms.

469 The detected water density varies from $5 \times 10^9 \text{ cm}^{-3}$ to 10^{13} cm^{-3} . The upper level of detected
470 water density is located at 40–60 km in the aphelion season and 70–90 km in the perihelion season.
471 At $L_s = 30\text{--}90^\circ$, during the northern summer, the maximal water density was observed in the
472 middle and high northern latitudes, corresponding to sublimation and subsequent release of vapor
473 from the northern polar cap. Most of the aphelion detections ($L_s = 30\text{--}120^\circ$) are located above the
474 hygropause with H_2O mixing ratio < 20 ppmv. Nevertheless, the hygropause was detected in the
475 middle northern latitudes at $L_s = 30\text{--}120^\circ$, manifested with a sharp decrease of the mixing ratio
476 from 50 to 5 ppmv, at the level of 15 km at 60°N , to 25–30 km at 40°N . Around the southern
477 equinox, in the $L_s = 150\text{--}210^\circ$ timeframe, low water amounts are distributed symmetrically
478 between high latitudes of both hemispheres.

479 The perihelion season is characterized by an increase of the water amount after $L_s = 240^\circ$ to
480 solstice, when the density reaches 10^{12} cm^{-3} and mixing ratio 100 ppm above 40 km. Water extends
481 to 90 km in the southern hemisphere confirming a seasonal impact on the hydrogen escape rate
482 (Bhattacharya et al., 2015). In the northern hemisphere water reaches 60–80 km and the solstice
483 pattern looks asymmetric. The largest mixing ratio of > 100 ppmv at 40–90 km was observed at
484 $L_s = 240\text{--}300^\circ$ in high and middle southern latitudes. At the end of the dusty season, during the
485 regional dust activity at $L_s = 310\text{--}330^\circ$, which occurs at this L_s in almost every year (Kass et al.,
486 2016), an abrupt increase of water was observed in both hemispheres in middle-to-high latitudes.

487 Intercomparison of water during the two global dust storms of MY28 and 34 shows that during
488 the MY28 GDS, coincident with summer solstice ($L_s = 268^\circ\text{--}285^\circ$) SPICAM observed a two- to
489 three-fold increase of the H_2O density at 60–80 km, exceeding that in regular years and in MY34
490 (when this period followed a GDS), especially in the northern hemisphere. The water density
491 increase during the MY34 GDS is comparable with the regular solstice water in the southern
492 hemisphere. In the northern hemisphere, even though our data for MY34 sample very high
493 latitudes only, the increase of water is prominent, marking an intensified circulation during global
494 dust events. We conclude that the MY28 GDS was exceptional in terms of delivering water to high

495 altitudes (80 km), exceeding by at least a factor of two the perihelion water maximum and the
496 water supply during the MY34 GDS.

497 Comparison of SPICAM/MEX and ACS NIR/TGO observations during the MY34 GDS shows
498 a good quantitative agreement and similar trends in both hemispheres.

499 The seasonal middle atmosphere water distribution presented here will help us to understand
500 in detail the water transport between hemispheres, to study its interannual variations, and to
501 separate the effect of seasonal water variation from the global and regional storm contribution, in
502 order to quantify better the hydrogen escape from Mars.

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512 **Competing interests**

513 The authors declare no competing interests

514 **Data and materials availability**

515 SPICAM raw and calibrated data are available in the ESA Planetary Science Archives (PSA)
516 <https://archives.esac.esa.int/psa/#!/Table%20View/SPICAM=instrument>. The H₂O density and

517 volume mixing ratio retrieved from the SPICAM occultation measurements are available at
518 <http://dx.doi.org/10.17632/vx4gks6bx7.1> (Fedorova, A. 2020b). The full IR Solar Occultation
519 Level-2B data will also be publicly available on the ESA PSA.

520

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